

COMPARISON BETWEEN COHESIVE ZONE MODELS AND A COUPLED CRITERION FOR PREDICTION OF EDGE DEBONDING

T. Vandellos^{a*}, E. Martin^b, D. Leguillon^c

^aONERA – The French Aerospace Lab, F-92322 Châtillon, France

^bLCTS, CNRS UMR 5801, Université de Bordeaux, 3 rue La Boétie, 33600 PESSAC, France

^cIJLRA, CNRS UMR 7190, Université Pierre et Marie Curie, 4 place Jussieu, 75252 PARIS Cedex 05, France

*thomas.vandellos@onera.fr

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Abstract

The onset of edge debonding within a bonded specimen submitted to bending is modeled with two numerical approaches: the coupled criterion and the cohesive zone model. The comparison of the results obtained with the both approaches evidences that (i) the prediction of edge debonding strongly depends on the shape of the cohesive law and (ii) the trapezoidal cohesive law is the most relevant model to predict the edge debonding as compared with the coupled criterion.

1. Introduction

The prediction of the onset of interface debonding is generally performed using a stress criterion [1] or linear fracture mechanics [2]. In each case, a characteristic length is needed which has to be identified experimentally. A first alternative approach is the use of a cohesive zone model [3] which simulates a progressive debonding build up in terms of continuum damage variables. It was shown by previous authors that the shape of the cohesive zone model (for a given value of the fracture toughness) does not have any influence on the steady state propagation of a rectilinear crack [4]. This result does not hold at crack initiation for which the shape of the cohesive law has a strong influence [4]. A second alternative approach is the coupled strength and energy criterion [5] which permits the prediction of the applied load involving the onset and the associated crack nucleation length. This approach has proved to be successful to analyze the onset of fracture mechanisms within composite materials [6,7,8] and bonded specimens [9,10]. An extension to the 3D geometry has been recently proposed [11] but, at the moment, the use of cohesive zone models seems to be easier.

The aim of this paper is to assess the capabilities of cohesive zone models to predict the failure onset of a bonded specimen submitted to bending as compared with the coupled criterion. In this study, it is important to note that the coupled criterion is just considered as a numerical reference solution; but this assumption should be experimentally verified. First, the both approaches are described. Second, several shapes of the cohesive law are compared to the coupled criterion in order to show their influence on the prediction. Finally, with the most relevant shape, the influence of the interfacial properties on the prediction of the edge debonding is presented to conclude about the use of the cohesive zone model.

2. The coupled criterion

As proposed by a previous author [5], combining an energy and a stress condition allows to derive an initiation criterion in the vicinity of a stress concentration. First, an energy balance between an elastic state prior to any crack growth and after the onset of a crack extension of area δS leads to the following incremental energy condition:

$$G_{inc}(a) = \frac{W(0) - W(a)}{\delta S} \geq G_C \quad (1)$$

where $W(0)$ is the potential energy at the initial state (without crack), $W(a)$ is the potential energy at the final state (with a crack of length a) and G_{inc} is the incremental energy release rate in which the infinitesimal energy rates of the classical Griffith approach are replaced by finite energy increments.

Second, a stress condition states that the normal out-of-plane stress σ along the anticipated path of crack nucleation is greater than the relevant strength σ_C

$$\sigma(x) \geq \sigma_C \quad \text{for } x \leq a \quad (2)$$

Finally, for a monotonic and increasing applied loading, the crack increment and the applied load at nucleation are obtained by combining the equations (1) and (2).

3. The cohesive zone models

Cohesive zone models are used to describe the behavior of interfaces. More precisely, the traction in mode I (*i.e.* the opening mode of fracture), T_1 (resp. T_2 in mode II (shearing mode)), between the top and bottom surfaces of the interface is related to the relative displacement in mode I δ_1 (resp. δ_2 in mode II). Several shapes of cohesive law (bilinear, trilinear and trapezoidal [12]), including an elastic part, can be obtained using the constitutive law written as

$$\begin{cases} \text{if } \delta_1 \geq 0 & \text{if } \delta_1 < 0 \\ \begin{cases} T_1 = K\delta_1(1 - \lambda) \\ T_2 = K\delta_2(1 - \lambda) \end{cases} & \begin{cases} T_1 = \alpha_C K\delta_1 \\ T_2 = K\delta_2(1 - \lambda) \end{cases} \end{cases} \quad (3)$$

where K is the initial stiffness of the interface, α_C is a penalization factor for out-of-plane compression, λ is the damage variable, related to the damage kinetics, and $f(\lambda)$ represents the effect of damage. The evolution of the damage variable λ is defined by

$$\left\{ \begin{array}{ll} \lambda = 0 & \text{if } \delta \leq \delta_0 \\ \lambda = (\delta_0 - \delta) \frac{\delta^* - \alpha_\sigma \delta_0}{\delta(\delta_0 - \delta^*)} & \text{if } \delta_0 < \delta \leq \delta^* \quad \text{and} \quad \dot{\lambda} \geq 0 \\ \lambda = \min \left\{ \frac{\delta(\delta^* - \delta_f) + \alpha_\sigma \delta_0(\delta_f - \delta)}{\delta(\delta^* - \delta_f)}, 1 \right\} & \text{if } \delta \geq \delta^* \end{array} \right. \quad (4)$$

The relative displacement δ is determined by

$$\delta = \sqrt{\langle \delta_1 \rangle_+^2 + (\delta_2)^2} \quad (5)$$

where $\langle \cdot \rangle_+$ are the classical Macaulay brackets defined by

$$\langle x \rangle_+ = \max\{0, x\} \quad (6)$$

It should be noted that, in order to avoid damage under pure out-of-plane compressive normal stress, the normal relative displacement δ_1 is only taken into account when positive. Parameters δ_0 and δ_f are material constants corresponding respectively to the relative displacement associated with the damage threshold σ_C and the interfacial stiffness K , and to the relative displacement attained when the energy release rate G is equal to the fracture toughness G_C . The threshold δ^* is the relative displacement associated with the maximal damageable stress σ^* which represents the admissible stress on the interface at the end of the first part of the damage process. It is defined by $\sigma^* = \alpha_\sigma \sigma_C$ where α_σ is a shape parameter. The value of δ^* depends on the shape parameter α_δ . The couple $(\alpha_\delta, \alpha_\sigma)$ defines where the negative slope of the cohesive law can be changed. Finally, the relative displacements δ_0 , δ^* and δ_f , represented in Figure 1, are defined by

$$\begin{cases} \delta_0 = \sigma_C / K \\ \delta^* = \delta_0 + \alpha_\delta (\delta_f - \delta_0) \quad \text{with } \alpha_\delta = \frac{G_C - 0.5\sigma_C \delta_0}{0.5\sigma_C (\delta_f - \delta_0)} - \alpha_\sigma \\ \delta_f = \frac{G_C - 0.5\sigma_C \delta_0}{0.5\sigma_C (\alpha_\delta + \alpha_\sigma)} + \delta_0 \end{cases} \quad (7)$$

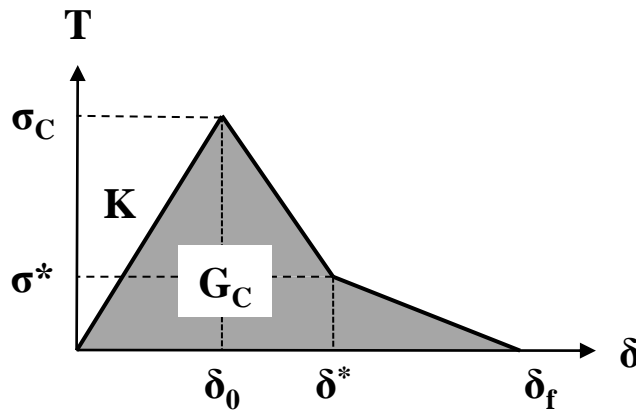


Figure 1. Representation of the relative displacements of the constitutive law from a tri-linear cohesive model

The shape of the cohesive law only depends on the values of the both shape parameters $(\alpha_\delta, \alpha_\sigma)$ which are equal to $(0,1)$ for the bilinear law and $(\alpha_\delta, 1)$ for the trapezoidal one. The three laws are illustrated in Figure 2 for identical interfacial properties (σ_C, G_C) . In the next paragraphs, the value of the relative displacement δ_0 has been imposed to 1.10^{-5} mm for all simulations. It is important to note that the parameter δ_0 is just considered as a numerical parameter. Moreover, two tri-linear shapes have been used, with $(\alpha_\delta, \alpha_\sigma) = (0.05, 0.9)$ and

$(\alpha_\delta, \alpha_\sigma) = (0.3, 0.6)$, and the shape parameter of the trapezoidal law has been imposed to $\alpha_\delta = 0.9$.

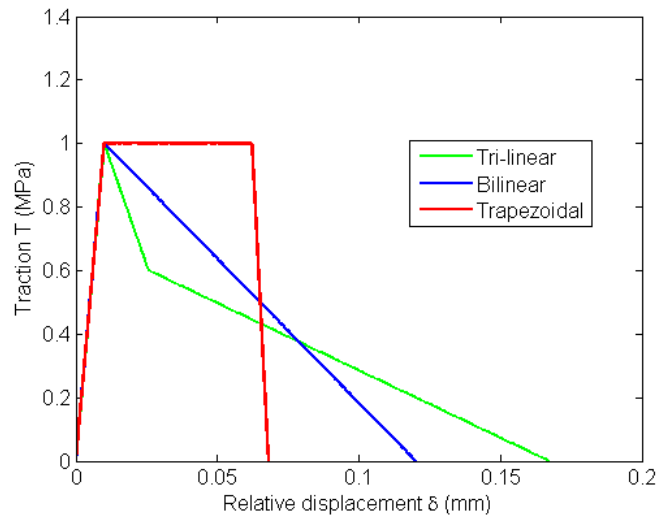


Figure 2. Illustration of the bilinear, tri-linear and trapezoidal shape of the cohesive zone model for the same interfacial properties

4. Prediction of edge debonding: coupled criterion versus cohesive zone models

In order to compare the two numerical approaches, simulations of a bonded specimen submitted to four-point flexure loading have been performed. Cohesive zone models and the coupled criterion are here applied to analyze the initiation of fracture mechanisms near the free edge between the bond and the substrate. The geometry of the specimen is schematized in Figure 3. It consists of two substrates with the same thickness $h = 2 \text{ mm}$ bonded with a thin interlayer which is here neglected. When the cohesive models are used, interface elements are inserted between the two bonded substrates. The elastic properties of the substrates are selected to be $E_S = 400 \text{ GPa}$ (Young's modulus) and $\nu_S = 0.2$ (Poisson's ratio). Due to the symmetry of the loaded specimen, it is assumed that the onset occurs near each free edge. A bidimensional finite element procedure with strongly refined mesh is used to derive the results which are now presented.

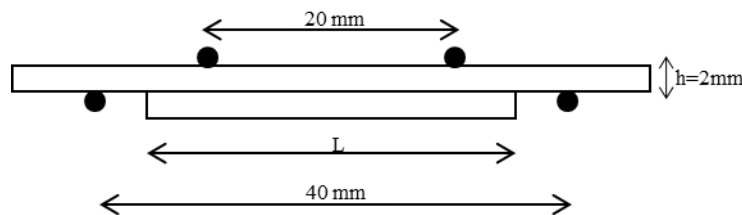


Figure 3. The geometry of the specimen tested under four-point bending

Several values of the interfacial strength σ_C and of the fracture toughness G_C have been studied for the comparison. These interfacial properties are considered similar whatever the fracture mode. The comparison between the load versus displacement curves obtained with the coupled criterion and with the cohesive zone models, for $\sigma_C = 1 \text{ MPa}$, (Figure 4) shows that the trapezoidal law is the most relevant model to predict the onset in a similar manner to

the coupled criterion, whatever the value of the fracture toughness. This observation is also true when $\sigma_C = 10$ MPa.

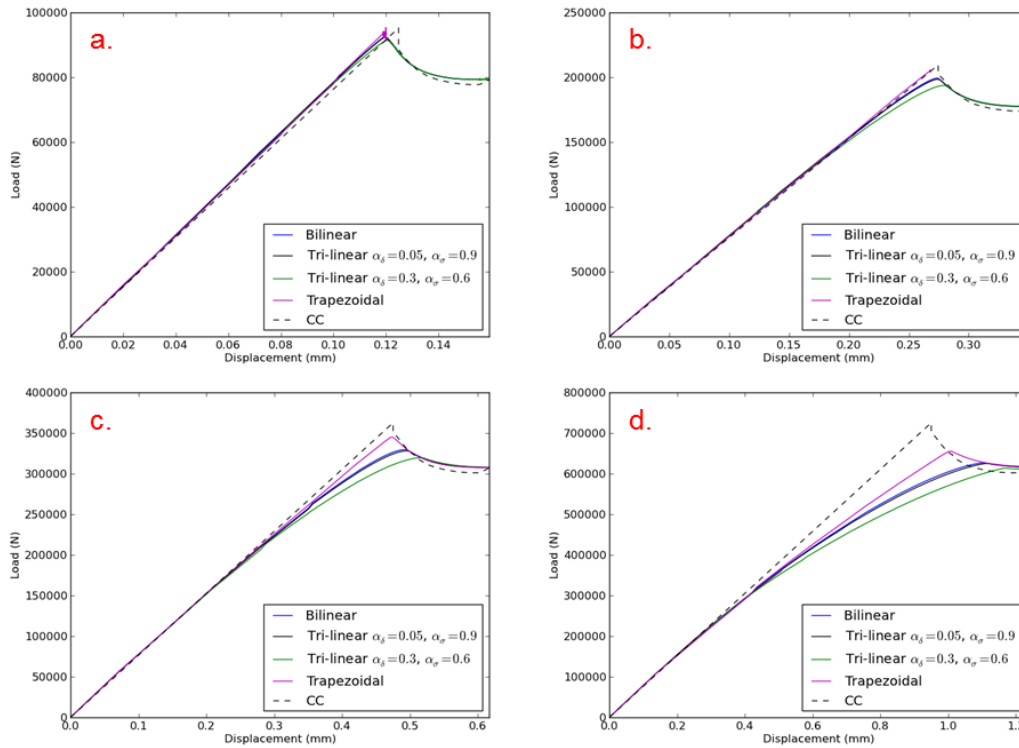


Figure 4. Comparison of load versus displacement curves between the coupled criterion (CC) and different cohesive zone models (CZM) with $\sigma_C = 1$ MPa and (a.) $G_C = 1$ J/m², (b.) $G_C = 5$ J/m², (c.) $G_C = 15$ J/m², (d.) $G_C = 60$ J/m²

This result can be explained comparing the damage kinetics of the different cohesive laws. Indeed, contrary to the bilinear and the tri-linear laws where the evolution of the damage variable λ between 0 (unbroken state) and 1 (broken state) is continuous and relatively slow, the damage variable of the trapezoidal model can evolve very quickly when the relative displacement δ exceeds δ^* (*i.e.* at the end of the plateau). This kinetics involves that the process zone (*i.e.* the area where the damage variable is positive but smaller than 1) with the trapezoidal law is smaller than the ones obtained with the bilinear and the tri-linear shapes, as shown in Figure 5. This observation could explain the better correlation obtained with the coupled criterion.

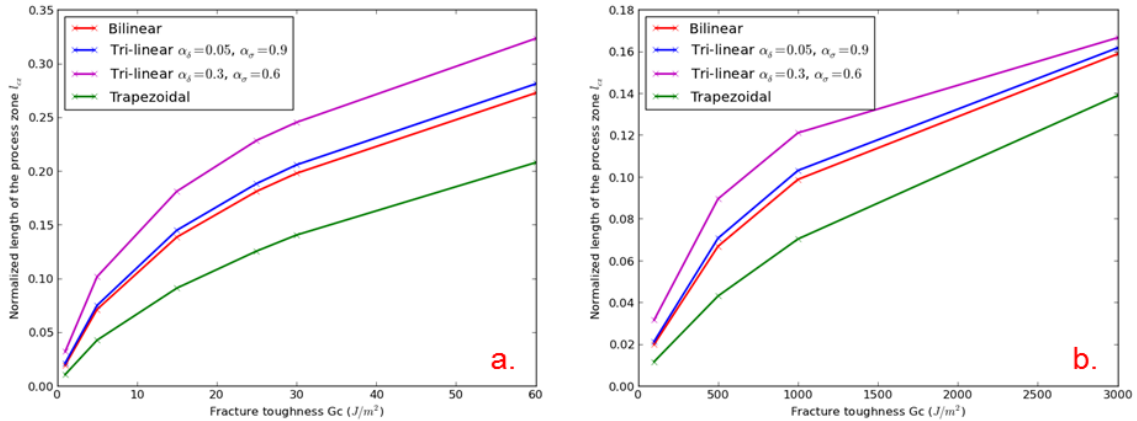


Figure 5. Normalized length of the process zone l_{cz}/L versus the fracture toughness G_C (a.) with $\sigma_C = 1$ MPa and (b.) with $\sigma_C = 10$ MPa.

Nevertheless, it is important to note that this correlation reduces when the fracture toughness increases (Figure 6 and Figure 7). This phenomenon results from the variation of the process zone length l_{cz} . Indeed, for the same interfacial strength σ_C , the larger the fracture toughness, the longer the length. Thus, the softening behavior becomes noticeable when the normalized crack length l_{cz}/L is higher than 10%. Consequently, as indicated in Table 1 and Table 2, the percentages error of the fracture displacement d_c and of the fracture load F_c exceed 5% as the consequence of softening (*i.e.* when $G_C > 15$ J/m² for $\sigma_C = 1$ MPa and when $G_C > 1000$ J/m² for $\sigma_C = 10$ MPa).

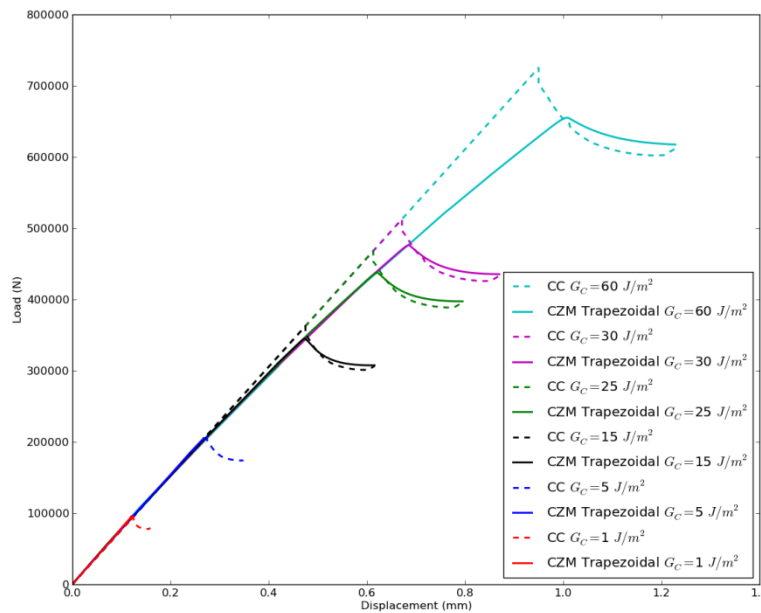


Figure 6. Comparison of load versus displacement curves between the coupled criterion (CC) and the trapezoidal cohesive zone model (CZM) for several fracture toughness with $\sigma_C = 1$ MPa

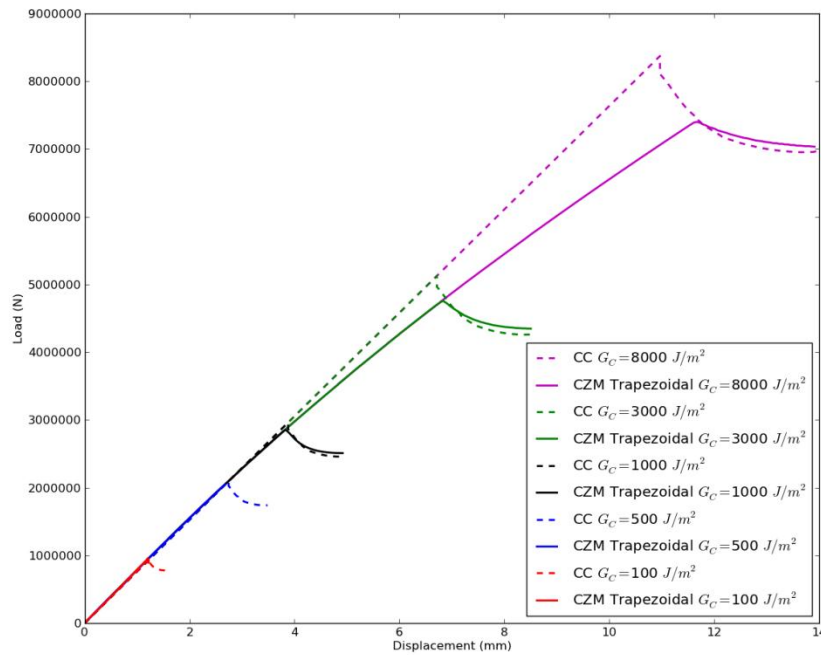


Figure 7. Comparison of load versus displacement curves between the coupled criterion (CC) and the trapezoidal cohesive zone model (CZM) for several fracture toughness with $\sigma_c = 10$ MPa

| G_c (J/m ²) | 1 | 5 | 15 | 25 | 30 | 60 |
|--|-----|-----|-------|------|------|------|
| $100 \left(\frac{d_c^{CC} - d_c^{CZM}}{d_c^{CC}} \right)$ | 4.7 | 2.4 | -0.08 | -1.5 | -2.7 | -6.3 |
| $100 \left(\frac{F_c^{CC} - F_c^{CZM}}{F_c^{CC}} \right)$ | 1.8 | 2.3 | 4.9 | 6.6 | 7.5 | 9.7 |

Table 1. Percentages error of the fracture displacement d_c and the fracture load F_c obtained by the trapezoidal cohesive zone model with $\sigma_c = 1$ MPa

| G_c (J/m ²) | 100 | 500 | 1000 | 3000 | 8000 |
|--|-----|-----|------|------|------|
| $100 \left(\frac{d_c^{CC} - d_c^{CZM}}{d_c^{CC}} \right)$ | 4.7 | 2.4 | 1 | -1.8 | -6.4 |
| $100 \left(\frac{F_c^{CC} - F_c^{CZM}}{F_c^{CC}} \right)$ | 1.8 | 2.2 | 3.8 | 7.2 | 11.6 |

Table 2. Percentages error of the fracture displacement d_c and the fracture load F_c obtained by the trapezoidal cohesive zone model with $\sigma_c = 10$ MPa

5. Conclusion

Simulations of a four-point bending test, for the prediction of edge debonding, with a coupled criterion and several cohesive zone models have been realized. First, the influence of the shape of the cohesive law on the prediction of the onset has been shown. The trapezoidal model appears the most adapted model to predict the initiation in a similar manner to the coupled criterion. Second, it has been demonstrated that the length of the process zone has to be small enough to verify a good correlation between the results obtained by the both

numerical approaches. Therefore, it seems possible to use the trapezoidal cohesive zone model, under a few material conditions, to predict precisely the edge debonding. A comparison between numerical and experimental results will be realized in future works in order to confirm the relevance of the trapezoidal cohesive zone model.

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